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# Computer Program for Three Body Break Up Reaction (Commemoration Issue Dedicated to Professor Sakae Shimizu on the Occasion of his Retirement)

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## Computer Program for Three Body Break Up Reaction

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A FORTRAN program is developed to calculate the energy spectrum for three body break up reaction at intermediate energy. Energy spectra are calculated on the basis of AGS (Alt, Grassberger and Sandhas) formalism assuming quasi free scattering with final state interaction. The calculated curves reproduced well energy spectra of  $^3\text{He}$  in coincidence with protons for  $^3\text{He}(\tau, \text{p})$  reaction and deuterons for  $^3\text{He}(\tau, \text{d})$  reaction at 120 MeV.

KEY WORDS Nuclear Reaction  $^3\text{He}(\tau, \text{p})$  at 120 MeV /  
Program for QF Scattering with FSI/

### I INTRODUCTION

Three body problems could be treated exactly on the basis of Faddeev's equations with given two body interaction potentials and the energy spectrum of nucleons from break up reaction of deuterons by nucleons could be successfully explained using a separable potential for nucleon-nucleon interaction<sup>1)</sup>. For the case where components of high angular momentum contribute to the two body interaction in the break up reaction, the problem is not so easy because that realistic separable potential has not yet been obtained and a long computer time needs to solve the integral equations. Furthermore if composite particles are concerned in the final state complicated reaction mechanism is presumed. At intermediate energies, however, continuous energy spectra of emitted particles show rather simple shape with large peaks due to quasi free scattering and some times with peaks for final state interaction appearing at low relative energies of the interacting two particles<sup>2)</sup>. Then it is reasonable to assume some simplifications in calculation for the break up reaction. Alt *et al.*<sup>3)</sup> have derived one dimensional integral equations for the three body scattering problem and derived a scattering amplitude at the special case where, in the first place, a quasi free process produces three particles and in the final state any two particles interact. We have obtained a computer program to calculate this amplitude and using it evaluated energy spectra for the break up reaction of  $^3\text{He}$  nucleus by  $^3\text{He}$  beam at 120 MeV.

### II EQUATIONS

Break up amplitude for three body break up reaction can be calculated using Eq. (5. 2) in Ref. 3 as follows,

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### Program for Three Body Break Up Reaction

$$X_{p_1, 1m}(z) = - \sum_{r \neq 1} \langle p_1 | T'_r(z) G_0(z) | 1, m; z \rangle + \sum_{r \neq 1} \langle p_1 | T_1(z) G_0(z) T_r(z) G_0(z) | 1, m; z \rangle, \quad (1)$$

where  $T'_r$  are non separable "amplitude" and small compared with the two body amplitude  $T_1$ . In the calculation the following conditions are assumed. 1) two body amplitude  $T_1$  and  $T'_r$  have not any singularity on the energy shell, 2) the Green's function  $G_0$  at the last position in the second term does not given any singularity because the cancellation of amplitudes at two integration angles  $p_1''$  and  $-p_1''$ , and 3) a pole approximation is applied to integrate the other Green's function in the second term. Then the cross section is given as follows,

$$\frac{d^3\sigma}{d\Omega_1 d\Omega_2 dE_1} = \frac{2\pi}{h^2} \cdot \frac{\mu_{in}}{k_{in}} \rho |M_+ + M_-|^2, \quad (2)$$

$$\begin{aligned} M_{\pm} &= \langle q_1 | X_{\pm} | q_0 \rangle \\ &= - \sum_{r \neq 1} \left( \frac{1}{S_{r1}} \right)^3 \langle p'_r | T'_r(z - q_1^2) | p_r \rangle \frac{1}{\{p'_1 - \frac{C_{r1}}{S_{r1}}(q_1 \mp q_0)\}^2 - Q} \\ &\quad \times \langle p'_1 - \frac{C_{r1}}{S_{r1}}(q_1 - q_0) | 1, m; z \rangle \\ &\quad + \sum_{r \neq 1} \left( \frac{1}{S_{r1}} \right)^3 \int d\Omega'' \langle p_1 | T_1(z - q_1^2) | p_1'' \rangle \frac{\pi}{2} i p_1'' \langle p'_r | T'_r(z - q_1^2) | p_r \rangle \\ &\quad \times \frac{1}{\{p'_1 - \frac{C_{r1}}{S_{r1}}(q_1 - q_0)\}^2 - Q} \langle p'_1 - \frac{C_{r1}}{S_{r1}}(q_1 \mp q_0) | 1, m; z \rangle, \end{aligned} \quad (3)$$

$$p_1''^2 = z - q_1^2,$$

where  $\rho$  is the phase space for factor three body system in the final state.  $M_+$  and  $M_-$  are matrix elements for the target break up process and projectile break up process, respectively. The variable  $p_1^2$  and  $q_1^2$  have a dimension of energy and relate to the usual wave vector  $k$  and mass  $m$ ,

$$\left. \begin{aligned} p_1 &= [2m_2 m_3 (m_2 + m_3)]^{-1/2} (m_3 k_2 - m_2 k_3), \\ q_1 &= [2m_1 (m_2 + m_3) (m_1 + m_2 + m_3)]^{-1/2} [m_1 (k_2 + k_3) - (m_2 + m_3) k_1], \end{aligned} \right\} \quad (4)$$

and the variables in the  $i$ - $j$ k system are related to the variables in the  $j$ - $ki$  system through the following relations,

$$\left. \begin{aligned} p_2 &= C_{21} p_1 + S_{21} q_1, \\ q_2 &= -S_{21} p_1 + C_{21} q_1. \end{aligned} \right\} \quad (5)$$

where

$$\left. \begin{aligned} C_{21} &= - \left[ \frac{m_1 m_2}{(m_1 + m_3) (m_2 + m_3)} \right]^{1/2} = C_{12}, \\ S_{21} &= \sqrt{1 - C_{21}^2}, \end{aligned} \right\} \quad (6)$$

Four types of potential are assumed in the calculation of the initial interaction, that is Gaussian type, Yukawa type, Woods Saxson type and Woods Saxson for real potential with the first derivative of Woods Saxson for imaginary potential,

$$\langle \mathbf{p}_r | T_r(z - q_r^2) | \mathbf{p}_r' \rangle = T_r(\mathbf{p}_r' - \mathbf{p}_r), \quad (7)$$

$$= -\pi^{3/2} \frac{V}{\alpha^3} \exp\left[-\frac{(\mathbf{p}_r' - \mathbf{p}_r)^2}{4\alpha^2}\right], \quad \text{for Gaussian} \quad (8)$$

$$= -4\pi V \left[ \frac{1}{\alpha_1^2 + (\mathbf{p}_r' - \mathbf{p}_r)^2} + \frac{\eta}{\alpha_2^2 + (\mathbf{p}_r' - \mathbf{p}_r)^2} \right], \quad \text{for Yukawa} \quad (9)$$

$$= \int d\mathbf{r}^3 V(r) e^{i(\mathbf{p}_r' - \mathbf{p}_r) \cdot \mathbf{r}}, \quad \text{for Woods Saxon } V(r). \quad (10)$$

The form factor of the target nucleus is assumed as,

$$\begin{aligned} \langle \mathbf{p}_1 | 1, m; z \rangle &= (z - q_0^2) \phi_{1m}(\mathbf{p}_1), \\ \phi_{1m}(\mathbf{p}_1) &= \frac{\sqrt{\alpha}}{\pi} \frac{1}{\alpha^2 + \mathbf{p}_1^2}. \end{aligned} \quad (11)$$

S wave final state interaction between particle 2 and particle 3 relates to the effective range parameters at zero energy limit in the two body scattering,

$$|\langle \mathbf{p}_1 | T_1(z - q_1^2) | \mathbf{p}_1' \rangle|^2 = \frac{C^2(\eta)}{4\pi^2} \cdot \frac{1}{C^2(\eta) q^2 + \left(-\frac{1}{a} + \frac{1}{2} r q^2\right)}, \quad (12)$$

$$\begin{aligned} q^2 &= (\mathbf{p}_1' - \mathbf{p}_1)^2, \\ C^2(\eta) &= 2\pi\eta / [\exp(2\pi\eta) - 1], \\ \eta &= z_2 z_3 e^2 / \hbar v, \end{aligned}$$

where  $C(\eta)$  is the Coulomb penetration factor. For the  $l$ -th partial wave interaction in the final state Brait-Wigner type resonance form is assumed as follows,

$$|\langle \mathbf{p}_1 | T_1(z - q_1^2) | \mathbf{p}_1' \rangle|^2 = A \frac{\Gamma_1^2}{(z - q_1^2 - E_1)^2 + \frac{1}{4} \Gamma_1^2} P_1(\cos\theta). \quad (13)$$

The angular integral of Eq. (3) is done analytically where the angular momentum of 1 is less than or equal to 2. The values of parameters in Eq. (8), (9), and (11) are transformed from the radial parameters in fm using following relation,

$$\alpha^2 \text{ MeV} = \frac{\gamma}{\mu a_{\text{fm}}}, \quad \gamma = \frac{41.6}{2},$$

$\mu$ : reduced mass

### III COMPUTER PROGRAM

The computer program "AGSIA" calculates the three body break up cross section on the basis of AGS formalism. Seven types of card have to read in the format of free field as the input data. Three body kinematics is calculated using a subprogram "KINEMA" which is the same as the program developed by Ohlsen<sup>4)</sup>, and the numerical integration of Eq. (10) is made using a sub-program "SIMPSON" which is a usual program of Simpson Integral. The matrix element in Eq. (3) is calculated using a sub-program "XMAT". Four types of initial interaction can be selected by **KCODE4** and five print out format can be selected by **KCODE3**. All calculated energy spectra are print out in units of  $10^{-26} \text{ cm}^2/\text{sr}^2 \cdot \text{MeV}$ . as a function of the laboratory energy of particle 1. The FORTRAN programs are given in an Appendix.

The energy spectrum of  $^3\text{He}$  in coincidence with protons for  $^3\text{He}(\tau, \tau p)d$  reaction

# Program for Three Body Break Up Reaction

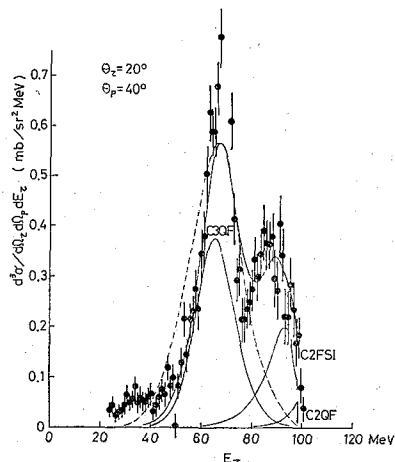


Fig. 1. Energy spectrum projected for the  ${}^3\text{He}(\tau, \tau p)d$  reactions at  $\theta_\tau=20^\circ$  and  $\theta_p=40^\circ$ . All solid curves are calculated on the basis of the AGS formalism. The dashed curve shows the calculated one of the simple spectator model.

at 120 MeV was calculated using the program of "AGSIA". Figure 1 shows the calculated curves at the angular set  $\theta_\tau=20^\circ$ ,  $\theta_p=40^\circ$  and  $\varphi_p=0^\circ$  (opposite side of, and in coplanar with, the beam) and shows the experimental energy spectrum which was obtained from the data in Ref. 2. The curves denoted as **C2QF** and **C2FSI** are calculated using only the first term at  $\gamma=2$  and the second term at  $\gamma=2$  in Eq. (1). The former corresponds to the cross section assuming the quasi free scattering of particle 1 ( ${}^3\text{He}$ ) and particle 3 (unobserved particle) and the latter corresponds to that for quasi scattering with the final state interaction of particle 2 and particle 3. The dashed curve is calculated using a simple impulsive spectator model. Nevertheless a normalization constant was needed to reproduce the experimental energy spectrum, the calculated energy spectrum is reproduced well using a reasonable parameter set.

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## REFERENCES

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## Appendix

## Instructions for use of "AGSIA"

## Card 1 (general conditions in the calculation)

:KCODE1, KCODE2, KCODE3, KCODE4, IK, NO, FACT

KCODE1 =1 print subtitle "XM3 IS NOT A COMPOSITE PARTICLE"

2 print subtitle "XM3=XM4+XM5 AND BREAKS TO SINGLET DEUTERON"

3 print subtitle "XM3=XM4+XM5 AND BREAKS TO TRIPLET p+n"  
SYSTEM"

if any number is chosen, the program does not change to calculate the reaction

KCODE2 =0 calculate the target breakup reaction;  $M=M_+$ 1 calculate the projectile breakup reaction;  $M=M_-$ 2 calculate the target breakup plus the projectile breakup reaction;  $M=M_+ + M_-$ 

KCODE3 =0 print out the kinematics of three body breakup reaction

1 print out all reaction amplitudes and phase space factor

2 print out the final cross sections

3 print out the cross sections for (1-2) quasi free, the cross sections for (1-3) quasi free and the final cross sections

4 print out all cross sections of components in Eq. (3) and the final cross sections

KCODE4 =0 initial interaction of Gaussian type

1 initial interaction of double Yukawa type

2 initial interaction of Woods-Saxson type

3 initial interaction of Woods-Saxson (real) with the first derivative of the Woods-Saxson (imaginary)

LK number of resonance in the final state interaction

NO number of angular set

FACT factor of the second term by the first term in Eq. (3), and set 1 in the usual calculation.

## Card 2 (kinematical parameters)

:XMP, XMT, XM1, XM2, XM3, Q, EPL, DE1L, ZZ

XMP mass of the projectile in atomic mass unit

XMT Mass of the target in atomic mass unit

XM1 Mass of the particle 1 in atomic mass unit

XM2 mass of the particle 2 in atomic mass unit

XM3 mass of the particle 3 in atomic mass unit

Q reaction Q value in MeV

EPL incident energy in MeV

DE1L energy increment in MeV

ZZ product of electric charges  $Z_2$  and  $Z_3$ 

## Card 3 (form factor and effective range parameters)

:ALPHTO, ALPHAO, ERANGT

ALPHTO width of the form factor in fm

ALPHAO scattering length in fm

ERANGT effective range in fm

## Card 4 (potential parameters for the first step interaction between particle 1 and particle 2.)

:VTPO1, RHTPO1, DHTPO1, YETAP1, VTPO2, RHTPO2, DHTPO2, YETAP2

VTPO1 real well depth in MeV

RHTPO1 real well width in fm

DHTPO1 real well diffuseness (for Gaussian or Woods-Saxson) or second well width in fm (for double Yukawa)

YETAP1 ratio of second Yukawa potential to the first Yukawa potential

VTPO2 imaginary well depth in MeV

RHTPO2 imaginary well width in fm

DHTPO2 imaginary well diffuseness (for Gaussian or Woods-Saxson) or the second well

# Program for Three Body Break Up Reaction

width in fm (for double Yukawa)

YETAP2 ratio of the imaginary depth of second Yukawa to that of first Yukawa

Card 5 (potential parameters for the first step interaction between particle 1 and particle 3)  
:VTDO1, RHTDO1, DHTDO1, YETAD1, VTDO2, RHXDO2, DHTDO2, YETAD2  
similar to the Card 4

Card 6 (resonance parameters for the S wave final state interaction)

:COEF(1), ER(1) WIDTH(1)

COEF(1) coefficient of the Brait-Wigner resonance

ER(1) resonance energy in MeV

WIDTH(1) resonance width in MeV

If the resonances for the P wave and the D wave final state interactions are calculated, the Card 7 and Card 8 are used for the resonance parameters. If any resonance is not calculate, the COEF(1) in the Card 6 must be equal to zero. And the Card of angular set follows to the resonance Card.  
Card 7 (angular set)

:TH1R, TH2R, PHI2

TH1R polar angle of particle 1 in degrees in the laboratory

TH2R polar angle of particle 2 in degrees

PHI2 azimuthal angle of particle 2 in degrees. Set zero when the detectors for particle 1 and particle 2 are set opposite side of, in coplaner, with, the incident beam.





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( 66 )

[illegible]

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0059 224 E2LH1N4C24C1C=1739922
      MAIN LOOP
0060 228 DO 10N L=1,300
0061 F=L
0062 IF (C1)=MHELL+WE
0063 IF (C1L)=L+M(WX) 103-101,131
0064 101 E1C=SUMH(1,2)=0+SUMH(ELL1,3)+A1+C=H1+A1+*2
0065 IF (C1)=SUMH(1,2)=L+C1L+M1(ELL1,3)+SUMH(1,2)C
0066 F123=L1C4M(WX)23
0067 F23(L)=L1U-1E23
0068 IF (C1)=L1U-10N-109-101
0069 102 W52=0+*W52M3=0+173/(EXP(WX*23)
0070 W52=0+*W52M3=0+173/(EXP(WX*23)
0071 W23=SUMH(2,2)=M1(ELL1,3)+M(WX)23
0072 W=0+*W23=H1C
0073 USUMX=M1(WX)12
0074 APM=H1M+5L1C(SUMH(2,2)+U+C2L
0075 RET=USUMX-U1U
0076 USUM=5L1+SUMH(2,2)=173/(EXP(WX*23)
0077 U2M=C1L+SUMH(2,2)=173/(EXP(WX*23)
0078 ALPHAM=H1M+S2L+C1M12L+U2M+C2L
0079 RET=U2M+USUM+U1U-173/(EXP(WX*23)
0080 F1=SUMH(2,2)=U+M1(ELL1,3)
0081 PC3(L)=H1+5L
0082 PCV3(L)=0+L
0083 PCV3(L)=0+L
0084 IF (C1)=H1+C1L+M1(ELL1,3)+M(WX)23
0085 PC(L)=SUMH(2,2)=M1(ELL1,3)+M(WX)23
0086 IF (ALPHAM+C2L-BETA) 109-230,103
0087 230 IF (RTA) 103,104,104
0088 103 H=1
0089 DO T=105
0090 104 H=1
0091 F2L(L,2)=0-0
0092 F1L(L,2)=0-0
0093 F12(L,2)=0+0
0094 TH2(L,2)=0-0
0095 H3C(L,2)=0-0
0096 PMQ3(L,2)=0-0
0097 PMQ3(L,2)=0-0
0098 TH2C(L,2)=0-0
0099 TH2C(L,2)=0-0
0100 TH2C(L,2)=0-0
0101 105 TH2C(L,2)=0
0102 PC1=L+L+L
0103 PTAB=SUMH(1,2)+PHAM+*BETA
0104 S103=C=1
0105 L1(L,1)=*L+*M(WX)1+ALPHAM+SUMH(2,2)
0106 P2=SUMH(2,2)=M1(ELL1,3)+M(WX)23
0107 P2M=H1+M1(ELL1,3)+M(WX)23
0108 P10M=W1+*P2=C1L+C2L+5L1+S2L+C1M12L
0109 DENH=SUMH(2,2)=M1(ELL1,3)+M(WX)23+*23
0110 RHQ3(L)=*M(WX)1+*M(WX)23+*M(WX)23
0111 DE2+*M(WX)23+ALPHAM+SUMH(2,2)=*M(WX)23+*M(WX)23
0112 DENH=SUMH(2,2)=M1(ELL1,3)+M(WX)23
0113 DENH=SUMH(2,2)=M1(ELL1,3)+M(WX)23
0114 TH2C(L,2)=0-0
0115 TH2C(L,2)=0-0
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0170 TH2C(L,2)=0-0
0171 TH2C(L,2)=0-0
0172 TH2C(L,2)=0-0
0173 TH2C(L,2)=0-0
0174 TH2C(L,2)=0-0
0175 TH2C(L,2)=0-0
0176 TH2C(L,2)=0-0
0177 TH2C(L,2)=0-0
0178 TH2C(L,2)=0-0
0179 TH2C(L,2)=0-0
0180 TH2C(L,2)=0-0
0181 TH2C(L,2)=0-0
0182 TH2C(L,2)=0-0
0183 TH2C(L,2)=0-0
0184 TH2C(L,2)=0-0
0185 TH2C(L,2)=0-0
0186 TH2C(L,2)=0-0
0187 TH2C(L,2)=0-0
0188 TH2C(L,2)=0-0
0189 TH2C(L,2)=0-0
0190 TH2C(L,2)=0-0
0191 TH2C(L,2)=0-0
0192 TH2C(L,2)=0-0
0193 TH2C(L,2)=0-0
0194 TH2C(L,2)=0-0
0195 TH2C(L,2)=0-0
0196 TH2C(L,2)=0-0
0197 TH2C(L,2)=0-0
0198 TH2C(L,2)=0-0
0199 TH2C(L,2)=0-0
0200 TH2C(L,2)=0-0

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K. FUJINAGA and S. KAKIJI

K. FUKUNAGA and S. KAKIGI

FAGS07370  
AGS07380  
AGS07390  
\*AGS07400  
AGS07410  
AGS07420  
\*AGS07430  
9FAGS07440  
AGS07450  
AGS07460  
AGS07470  
AGS07480  
AGS07490  
AGS07500  
AGS07510